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RABI OSCILLATION IN COUPLED SYSTEM OF COOPER PAIR BOX AND CAVITY PHOTON

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The quantum coherent oscillations of a Josephson charge qubit prepared in a superposition of two charge states. A purpose of different quantum computing scheme based on Josephson charge qubits. Multiphoton processes was studied in Josephson charge qubit in contrast to the usual Jaynes-Cumming model, the Hamiltonian includes higher order interactions between the two levels system and the non-classical microwave field. M.A. Nielsen [1], demonstrated the coherent manipulation of the qubit state in Rabi oscillation. The temporal coherent oscillation of the population inversion in a two level system driven by external electromagnetic field is known as Rabi oscillation [2].

Introduction

Rabi oscillations are periodical transitions of a two state quantum system between its stationary states in the presence of an oscillatory driving field [3]. Theoretically predicted by Rabi in 1937 [4], Rabi oscillations were firstly observed by Torrey in 1949 on nuclear spins in radio frequency magnetic field [5]. Rabi oscillations were performed in many other systems, such as electromagnetically driven atoms [6], semiconductor quantum dots [7], Josephson qubits [8], spin qubits [9].

Mathematical calculation

The Hamiltonian for a spherical of Ta/In Ox/Ta quantum dotes in effective mass approximation is given by

$$H = -\nabla^2 + v l_z + v^2 \rho^2 / 4 - V_b(\mathbf{r}) - 2/r$$

The rms vacuum electric field amplitude E_{vac} in a mass of frequency is

$$E_{vac} = \left(\frac{\hbar\omega}{2\varepsilon_0 V}\right)^{1/2}$$

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where ε_0 is the permittivity of free space is in the size of an arbitrary quantization volume and $\hbar \omega$ are in S.I. unit. The coupling of an atom to a field mode is described by frequency

$$-\Omega_{ef} = \frac{D_{ef} - E_{vac}}{\hbar}$$

The Josephson phase qubit is simply a hysteric current biased Josephson junction that is well isolated from dissipation. The Hamiltonian is-

$$H = 4E_c \hbar^{-2} P_{\nu}^2 - E_j \left\{ \cos \nu + \left(\frac{I}{I_c}\right) \nu \right\}$$

where v is Gauge invariant.

The experimental value for effective mass approximation given data for (figure) for Rabi frequency in table 1

S. No.	Metal	Atomic Mass	$\begin{aligned} \text{Rabi-frequency } (0 \rightarrow 1, \text{ transition}) \\ \Omega_{R,01}^{(1)} &= \Omega_{01}^{(1)} \times \\ & \left(1 - \frac{[\Omega_{01}^{(1)}]^2}{4(\omega_{01-}\omega_{12})}\right) \text{MHz} \end{aligned}$	Micro- wave current I _{ac}	The a.c. stark shift $\Delta \omega_{01}$ of the one photon
1	In	114.82	68	2.50	15
2	Al	269.82	96	3.75	20
3	As	799.22	110	8.75	40

Table-1

Author has calculated value for x = 1 given data for $(0 \rightarrow 1$ transition) Rabi frequency in table 2.

S. No.	Metal	Atomic mass	$\begin{aligned} & \text{Rabi-frequency} \\ & (0 \rightarrow 1, \text{transition}) \\ & \Omega_{R,01}^{(1)} = \Omega_{01}^{(1)} \times \\ & \left(1 - \frac{[\Omega_{01}^{(1)}]^2}{4(\omega_{01-}\omega_{12})}\right) \text{MHz} \end{aligned}$	Microwave current I _{ac}	The a.c. stark shift $\Delta \omega_{01}$ of the one photon
1.	Та	180.95	66	2.50	14.98
2.	In	114.82	95	3.75	20.45
3.	0	15.99	108	8.75	40.35

Table- 2



Experimental value for Rabi frequency $(0 \rightarrow 1, \text{ transition})$ agrees well with the calculated value for *Ta/In Ox*. High power Rabi Oscillation can in principle be used for high fidelity microwave central of super conduction phase qubit. Multiphoton Rabi Oscillation and a.c. stark shift can also be explained.

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